Stochastic sampling for the analytic continuation of imaginary-time data

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$$G(\tau) = \pm \int_{-\infty}^{\infty} d\omega \, \frac{e^{-\tau\omega}}{1 \pm e^{-\beta\omega}} \, A(\omega)$$





outline

- analytic continuation
- singular value decomposition
- ill-posedness and role of non-negativity constraint
- average spectrum approach
- blocked mode sampling
- results depend on discretization
- separating numerics from default model
- a practical approach

Green/correlation function

spectral function $A(\omega)$ defines analytic function in upper/lower half of complex frequency plane

$$G(z) = \int_{-\infty}^{\infty} d\omega \, \frac{A(\omega)}{z - \omega}$$

reconstruct
$$A(\omega)$$
 from $G(i\omega_n)$ $\omega_n = \begin{cases} (2n+1)\pi/\beta & \text{fermions} \\ 2n\pi/\beta & \text{bosons} \end{cases}$
 $G(i\omega_n) = \int_{-\infty}^{\infty} d\omega \frac{A(\omega)}{i\omega_n - \omega}$

reconstruct $A(\omega)$ from $G(\tau)$

$$G(\tau) = \frac{1}{\beta} \sum_{n} G(i\omega_{n}) e^{-i\omega_{n}\tau} = \pm \int_{-\infty}^{\infty} d\omega \frac{e^{-\tau\omega}}{1 \pm e^{-\beta\omega}} A(\omega)$$

analytic continuation

 $A(\omega) \ge 0$ for fermions

 $A(-\omega) = -A(\omega)$ for bosons: consider $A(\omega)/\omega \ge 0$ for $\omega \ge 0$

$$G(i\omega_n) = \int_0^\infty d\omega \, \frac{2\omega^2}{\omega_n^2 + \omega^2} \, \frac{A(\omega)}{\omega}$$

e.g. optical conductivity

$$\Pi(i\omega_n) = \frac{1}{\pi} \int_0^\infty d\omega \ \frac{\omega^2}{\omega_n^2 + \omega^2} \quad \sigma(\omega)$$

$$\Pi(\tau) = \frac{1}{\pi} \int_0^\infty d\omega \; \frac{\omega e^{-\tau\omega}}{1 - e^{-\beta\omega}} \; \sigma(\omega)$$

Fredholm equation

$$g(y) = \int K(y, x) f(x) \, dx$$

data kernel model

$$|g\rangle = K|f\rangle$$

inverse problem: given g(y) find f(x)

ill posed

singular value decomposition

$$|g\rangle = K|f\rangle$$

SVD of kernel

 $K = USV^{\dagger}$

 $V = (|v_1\rangle, |v_2\rangle, ...)$ orthonormal basis in model space $U = (|u_1\rangle, |u_2\rangle, ...)$ orthonormal basis in data space

S diagonal rectangular matrix with $s_i \ge s_{i+1}$ (singular values)

spectral representation of kernel

$$K|v_i\rangle = s_i|u_i\rangle$$

singular values and singular modes



least-squares solution

$$|g\rangle = K|f\rangle$$

to find best f, minimize $||K|f\rangle - |g\rangle||^2$

$$|f\rangle = \sum |v_i\rangle \langle v_i | f \rangle = \sum f_i |v_i\rangle \quad \rightsquigarrow \quad K|f\rangle = \sum f_i s_i |u_i\rangle$$
$$|g\rangle = \sum |u_i\rangle \langle u_i | g \rangle = \sum g_i |u_i\rangle$$

$$\frac{\partial}{\partial \bar{f_i}} \| K|f\rangle - |g\rangle \|^2 = s_i \left(f_i s_i - g_i\right) \stackrel{!}{=} 0$$

$$f_i = \frac{g_i}{s_i}$$

ill posedness

$$|f_{LS}\rangle = \sum \frac{g_i}{s_i} |v_i\rangle$$

modes with small singular value hardly affect data small errors in data are catastrophically amplified

the affected modes tend to be highly oscillating sawtooth noise in least-squares fit

Picard condition
$$\||f\rangle\|^2 = \sum |\langle v_i | f \rangle|^2 = \sum \frac{|\langle u_i | g \rangle|^2}{s_i^2} < \infty$$

implies smoothness of data noise in data leads to violation of Picard condition

 $f_i = g_i / s_i$ variance in data $\sigma_i \Rightarrow$ variance in model: σ_i / s_i

ill posedness



Tikhonov regularization

minimize
$$\|K|f\rangle - |g\rangle\|^2 + \alpha \||f\rangle\|^2$$

$$f_i = \frac{s_i^2}{s_i^2 + \alpha^2} \frac{g_i}{s_i}$$

determine regularization parameter a



effect of regularization



effect of non-negativity constraint



non-negativity constraint regularizes singular modes: modes with small *s_i* adjust such that leading modes can be optimized



Bayesian inference

take quality of data and a priori knowledge about model into account $|g\rangle = K |f\rangle$

data given by mean $ar{g}$ and covariance Σ $P(g|ar{g},\Sigma)\propto e^{-rac{1}{2}\langle g-ar{g}|\Sigma^{-1}|g-ar{g} angle}$

probability of model f given measured data \bar{g} , Σ

$$P(f|\bar{g}, \Sigma) = \frac{P(\bar{g}, \Sigma|f) P(f)}{P(\bar{g}, \Sigma)}$$

likelihood
$$P(\bar{g}, \Sigma|f) = \int \mathcal{D}g \ P(\bar{g}|g) P(g|f) = P(\bar{g}, \Sigma|Kf)$$
$$\propto e^{-\frac{1}{2}\langle Kf - \bar{g}|\Sigma^{-1}|Kf - \bar{g}\rangle} \xrightarrow{\text{assume}}{P(g) = P(\bar{g}, \Sigma)}$$

prior probability *P(f)*: *a priori* information about model

likelihood function

$$P(\bar{g},\Sigma|f)\propto e^{-rac{1}{2}\langle Kf-ar{g}|\Sigma^{-1}|Kf-ar{g}
angle}$$

absorb covariance in data/kernel

Cholesky decomposition $\Sigma^{-1} = W^{\dagger}W$

$$\begin{split} \langle Kf - \bar{g} | \Sigma^{-1} | Kf - \bar{g} \rangle &= \langle Kf - \bar{g} | W^{\dagger} W | Kf - \bar{g} \rangle \\ &= \langle WKf - W\bar{g} | WKf - W\bar{g} \rangle \\ &= \langle \tilde{K}f - \tilde{g} | \tilde{K}f - \tilde{g} \rangle \end{split}$$

prior probabilities

• Tikhonov

$$P(f) \propto \exp\left(-\frac{\alpha}{2} \langle f|f\rangle\right)$$

favor solutions with small norm

• MaxEnt $P(f) \propto \exp\left(-\alpha \int f(x) \ln\left(f(x)/d(x)\right) dx\right)$

favor solutions close to default model d(x)

• non-informative

$$P(f) \propto \begin{cases} 0 & \text{if } \exists x : f(x) < 0 \\ 1 & \text{otherwise} \end{cases}$$

all a priori admissible solutions equally probably

how to pick model given $P(f | \bar{g}, \Sigma)$?

minimize cost of being wrong best estimate minimizes expected loss

$$\int \mathcal{D}f L(\bar{f}, f) P(f|\bar{g}, \Sigma)$$

loss function $L(\dot{f},f)$

 $L(\bar{f}, f) = \begin{cases} 0 & \text{when } \bar{f} = f \\ 1 & \text{otherwise} \end{cases} \quad \rightsquigarrow \quad \bar{f} \text{ maximizes } P(f|\bar{g}, \Sigma)$

$$L(\overline{f},f) = \||\overline{f} - f\rangle\|^2 \quad \rightsquigarrow \quad \overline{f} = \int \mathcal{D}f \, f \, P(f|\overline{g},\Sigma)$$

average over posterior

respects constraint if set of admissible functions is convex

Stochastic Sampling

The Average Spectrum Method for the Analytic Continuation of Imaginary-Time Data

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In this paper we present the average spectrum method, a new method for obtaining realfrequency information from imaginary-time quantum Monte Carlo data. This technique does not require the adjustable parameters, smoothness constraints, or model forms of some previous techniques, yet produces smooth, consistent spectra from noisy data. Various tests of the method on mock data are presented, as well as realistic applications to the two-dimensional Hubbard model.

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Stochastic Sampling

impose *a priori* via constraints variance in data smoothes model

• S.R. White:

The Average Spectrum Method for the Analytic Continuation of Imaginary-Time Data Springer Proceedings in Physics **53**, 145 (1991)

- A.W. Sandvik: Stochastic method for analytic continuation of quantum Monte Carlo data Phys. Rev. B 57, 10287 (1998)
- K. Vafayi and O. Gunnarsson: Analytical continuation of spectral data from imaginary time axis to real frequency axis using statistical sampling Phys. Rev. B 76, 035115 (2007)
- K.S.D. Beach:

Identifying the maximum entropy method as a special limit of stochastic analytic continuation cond-mat/0403055

• O.F. Syljuåsen:

Using the average spectrum method to extract dynamics from quantum Monte Carlo simulations

Phys. Rev. B 78, 174429 (2008)

Gaussian distribution in f

$$|\bar{f}
angle \propto \int_{\mathcal{F}} \mathcal{D}f \, e^{-rac{1}{2}\langle \tilde{K}f - \tilde{g}|\tilde{K}f - \tilde{g}
angle} |f
angle$$

rewrite as Gaussian in f

pseudo-inverse of kernel K

$$K = \sum_{i} |u_i\rangle s_i \langle v_i| \qquad K^- = \sum_{s_i>0} \frac{|u_i\rangle \langle v_i|}{s_i}$$

$$\langle \tilde{K}f - \tilde{g} | \tilde{K}f - \tilde{g} \rangle = \langle f - \tilde{K}^{-}\tilde{g} | \tilde{K}^{\dagger}\tilde{K} | f - \tilde{K}^{-}\tilde{g} \rangle - \langle g_{\perp} | g_{\perp} \rangle$$

$$\sum_{i} \tilde{s}_{i}^{2} |f_{i}|^{2} - \tilde{s}_{i}(f_{i}^{*}\tilde{g}_{i} + \tilde{g}_{i}^{*}f_{i}) + |\tilde{g}_{i}|^{2} = \sum_{\tilde{s}_{i}>0} \tilde{s}_{i}^{2} \left| f_{i} - \frac{\tilde{g}_{i}}{\tilde{s}_{i}} \right| + \sum_{\tilde{s}_{i}=0} |\tilde{g}_{i}|^{2}$$

 g_{\perp} : component of data that cannot be described by model

Gibbs sampling

write model in (finite) basis $|f\rangle = \begin{pmatrix} f_1 \\ f_2 \\ \vdots \end{pmatrix}$

sample component $f_i \rightarrow f_i$ keeping all other components fixed

$$|f\rangle = \begin{pmatrix} f_i \\ f_\perp \end{pmatrix} \quad \tilde{K} = (\tilde{K}_i, \tilde{K}_\perp) \quad \tilde{g}_i = \tilde{g} - \tilde{K}_\perp f_\perp$$

$$P(f_i \to f_i' | f_\perp) = \exp\left(-\frac{1}{2} \frac{\langle \tilde{K}_i f_i' - \tilde{g}_i | \tilde{K}_i f_i' - \tilde{g}_i \rangle}{\langle \tilde{K}_i f_i - \tilde{g}_i | \tilde{K}_i f_i - \tilde{g}_i \rangle}\right)$$

Gaussian with
$$\sigma^2 = 1/ ilde{K}_i^\dagger ilde{K}_i \quad \mu = \sigma^2 \, ilde{K}_i^\dagger ilde{g}_i$$

non-negativity constraint limits allowed values → truncated Gaussian

can be efficiently sampled

mode sampling

in SVD basis the singular modes have independent Gaussians

$$P(f_i \to f'_i | \mathbf{f}) = \exp\left(-\frac{1}{2} \frac{|f'_i - \tilde{g}_i / \tilde{s}_i|^2}{|f_i - \tilde{g}_i / \tilde{s}_i|^2}\right)$$

but are coupled by constraint $f(\omega) = \sum f_i v_i(\omega) \ge 0$

coupling of modes reduces variance of modes with small s_i

non-negativity constraint regularizes singular modes: modes with small *s_i* adjust such that leading modes can be optimized



component sampling

in frequency basis $f(\omega_i)$

non-negativity constraint: $f(\omega_i) \ge 0$

 $P(f(\omega_i) \rightarrow f'(\omega_i)|f_{\perp}) =$ Gaussian with

$$\sigma^{-2} = \int d\tau \, |\tilde{K}(\tau,\omega_i)|^2 \quad \mu = \sigma^2 \int d\tau \, g_i(\tau) \, K^*(\tau,\omega_i)$$

choice of basis

mode sampling



component sampling



best at constraint

best at maximum

efficiency: thermalization



efficiency: autocorrelation









blocked mode sampling

subdivide ω -axis into subintervals on which to do mode samping



use different primes for subdivision to avoid seams

starting point

NNLS solutions tend to be on boundary of admissible set take linear combination of several NNLS solutions obtained from adding noise to data



test cases

PHYSICAL REVIEW B 82, 165125 (2010)

Analytical continuation of imaginary axis data for optical conductivity

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We compare different methods for performing analytical continuation of spectral data from the imaginary time or frequency axis to the real frequency axis for the optical conductivity $\sigma(\omega)$. We compare the maximum entropy (MaxEnt), singular value decomposition (SVD), sampling, and Padé methods for analytical continuation. We also study two direct methods for obtaining $\sigma(0)$. For the MaxEnt approach we focus on a recent modification. The data are split up in batches, a separate MaxEnt calculation is done for each batch and the results are averaged. For the problems studied here, we find that typically the SVD, sampling, and modified MaxEnt methods give comparable accuracy while the Padé approximation is usually less reliable.

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optical conductivity
$$\Pi(\nu_n) = \frac{2}{\pi} \int_0^\infty \frac{\omega^2}{\nu_n^2 + \omega^2} \sigma(\omega) \, d\omega$$

$$\sigma(\omega) = \left\{ \frac{W_1}{1 + (\omega/\Gamma_1)^2} + \frac{W_2}{1 + [(\omega - \epsilon)/\Gamma_2]^2} + \frac{W_2}{1 + [(\omega + \epsilon)/\Gamma_2]^2} \right\} \frac{1}{1 + (\omega/\Gamma_3)^6}$$

$$\Gamma_1 = 0.3 \text{ or } 0.6, \Gamma_2 = 1.2, \Gamma_3 = 4, \epsilon = 3, W_1 = 0.3, W_2 = 0.2$$

vary noise on data

stochastic sampling



improve cutoff: $\omega_{cut} = 8$



improve cutoff: $\omega_{cut} = 16$



improve cutoff: $\omega_{cut} = 32$



cutoff dependence



same model, more accurate data

effect of discretization





$$ho(\omega) = rac{\gamma}{\pi} rac{1}{\omega^2 + \gamma^2}$$



behavior where data is weak



for large ω stochastic sampling result behaves like density function $\varrho(\omega)$

behavior where data is weak



for large ω stochastic sampling result behaves like density function $\varrho(\omega)$

limit of no data



in absence of data stochastic sampling recovers grid density default model

understand grid dependence?



grid resolution



Gamma distribution



 $\int_0^x \operatorname{Gamma}(k_1, \lambda; x') \operatorname{Gamma}(k_2, \lambda; x - x') \, dx = \operatorname{Gamma}(k_1 + k_2, \lambda; x)$

simulating grids of different resolution

exponential grid with 32 and 128 points



grid type



$$\frac{p_{\text{fine}}(f_1, f_2, \dots, f_n)}{p_{\text{other}}(f_1, f_2, \dots, f_n)} \propto \lim_{\lambda \to \infty} \frac{\prod_{i=1}^n f_i^{k_i - 1} e^{-f_i / \lambda}}{\prod_{i=1}^n e^{-f_i / \lambda}} = \prod_{i=1}^n f_i^{k_i - 1}$$

simulating grids of different type

exponential and Lorentzian grid



general case

simulate grid of n_1 points with density function $\varrho_1(x)$ on grid of n_2 points with density function $\varrho_2(x)$

reweight model $(f_1, f_2, ..., f_{n2})$ on second grid by

$$\frac{p_2(f_1, f_2, \dots, f_n)}{p_1(f_1, f_2, \dots, f_n)} = \prod_{i=1}^n f_i^{\frac{n_1 \rho_1(x_i)}{n_2 \rho_2(x_i)} - 1}$$

make result in absencence of data proportional to $\varrho_1(x)$

separate grid from default model

recipe for discretization

estimate width of model from NNLS

$$\mu = \int f(\omega) \, d\omega \quad \sigma^2 = \int (\omega - \mu)^2 f(\omega) \, d\omega$$

default model: least informative function that reproduces lowest moment

$$\omega \ge 0: \qquad D(\omega) = \frac{1}{\mu} e^{-\omega/mu}$$

otherwise:
$$D(\omega) = \frac{1}{\sqrt{2\pi}\sigma} e^{-\frac{(\omega-\mu)^2}{2\sigma^2}}$$

grid: make numerical error in evaluating Kernel smaller than noise in data

ω ≥ 0: Lorentzian at 0 with half-width µotherwise: Lorentzian at µ with width σ

does NNLS depend on grid?



does NNLS depend on grid?

different grids, same peaks (up to resolution)



results



reduce noise on data



second peak further out



summary



regularization by constraint



avrage spectrum method

$$|ar{f}
angle \propto \int_{\mathcal{F}} \mathcal{D}f \ e^{-rac{1}{2}\langle ilde{\kappa}f - ilde{g} | ilde{\kappa}f - ilde{g}
angle} \left|f
ight
angle$$

blocked mode sampling

separate grid & default model

